

Mechanics and Physics of Porous Solids

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Based on a lecture from O. Coussy and M. Vandamme



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- 5 The unsaturated porous solid
- 6 Confined phase transitions
- 7 Experimental considerations

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Definition

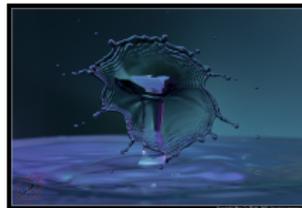
What is mechanics ?

Mechanics is the study of the movement and deformation of physical systems.

A porous solid consists in a solid matrix filled by one or more fluid (gas or liquid). To consider the mechanics of a porous solid we need to know the *mechanics of deformable solids* and *the mechanics of fluids* and the coupling between both.



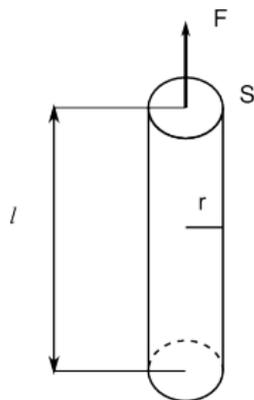
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Simple deformations, elasticity

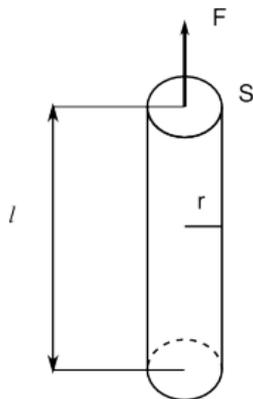


Cylinder subjected to a force F

The strain (deformation) of the cylinder is defined as: $\varepsilon = \frac{\Delta l}{l_0}$

The stress applying on the cylinder is defined as $\sigma = \frac{F}{S}$

Simple deformations, elasticity



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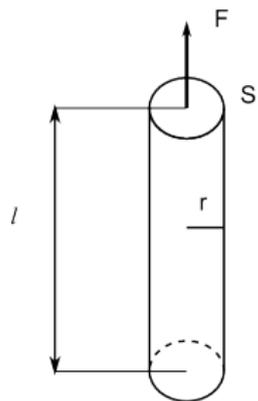
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Hooke's law: *ut tensio sic vis*

If the stress is moderate, the cylinder is within the elastic domain:

- the deformation is linear with the stress $\sigma = E\varepsilon$, E is called Young Modulus ($\sim GPa$)
- the deformation is reversible

Simple deformations, elasticity

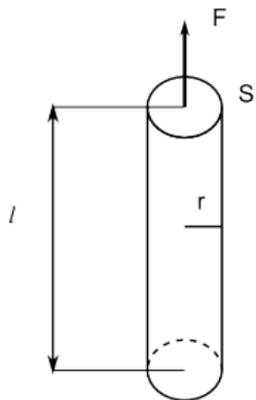


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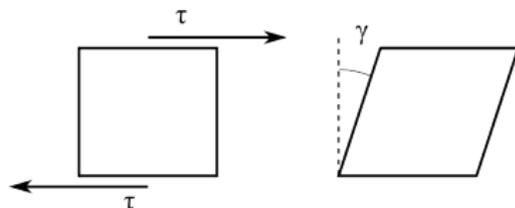
Poisson, bulk moduli

If the stress is moderate, the cylinder is within the elastic domain:

- the change of radius is characterized by the Poisson modulus: $\frac{\Delta r}{r_0} = -\nu \cdot \varepsilon$
- for an isotropic deformation (pressure), the volume change is: $\Delta P = -K \frac{\Delta V}{V}$

Simple deformations, elasticity

Shear



If the force is parallel to the surfaces it is called shear

The variation of the angle γ (shear angle) is also linear under the elasticity assumption

$$\tau = G\gamma, \text{ with } G \text{ the shear modulus}$$

Relation between the different moduli

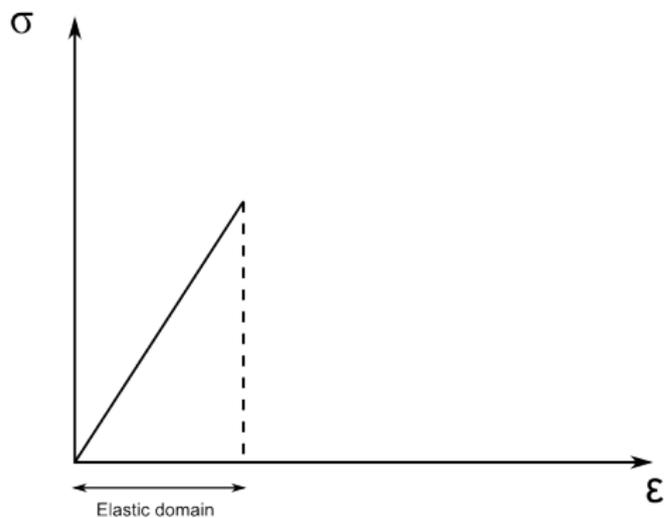
- $G = \frac{E}{2(1+\nu)}$
- $\frac{1}{E} = \frac{1}{9K} + \frac{1}{3G}$
- $K = \frac{1}{3} \frac{E}{1-2\nu}$
- if $\nu > 0.5$ the volume increases under isotropic compression \Rightarrow Non physical
- if $\nu < 0.5$ the volume decreases under compression and increases under traction

Usual values of the mechanical moduli for different materials

Material	E (GPa)	G (GPa)	ν	K (GPa)
Rubber	0.001 to 0.1	0.0003	0.48 to 0.5	0.0001 to 0.1
Steel	210	77	0.3	160
Sandstone	3 to 90	1 to 40	0.2 to 0.35	2 to 110
Limestone	9 to 80	3 to 30	0.2 to 0.3	5 to 70
Granite	10 to 70	4 to 30	0.1 to 0.2	50
Concrete	20 to 50	21	0.1 to 0.3	15
Cork	0.0186	0.0093	0	0.0062

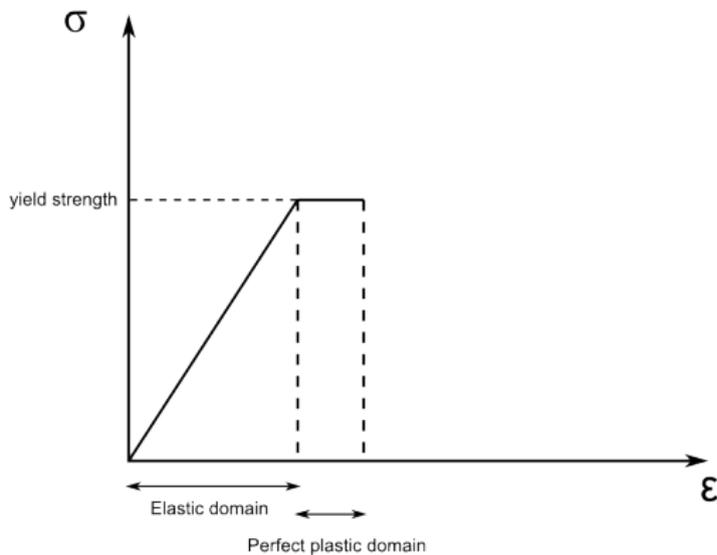
Beyond elasticity, plasticity and rupture

Response of a system under a simple traction test



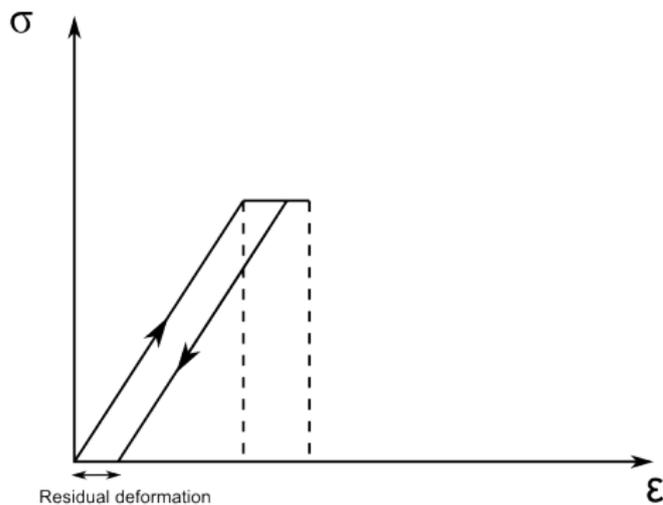
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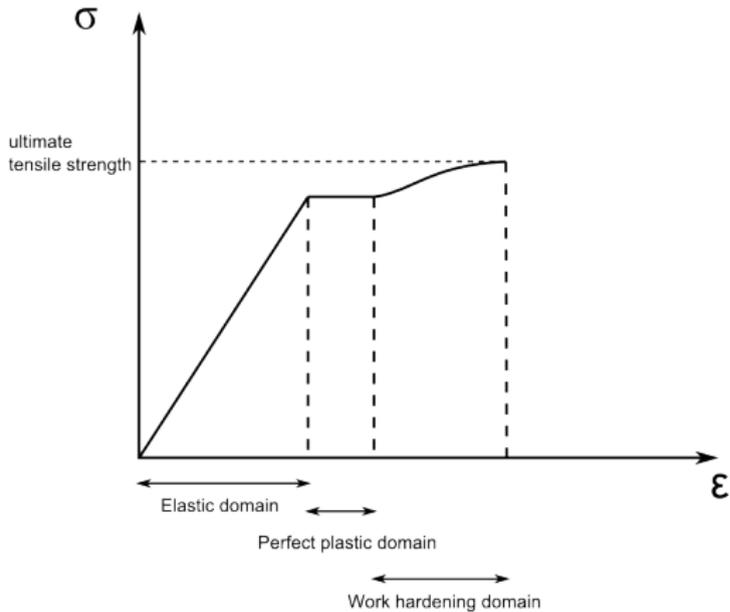
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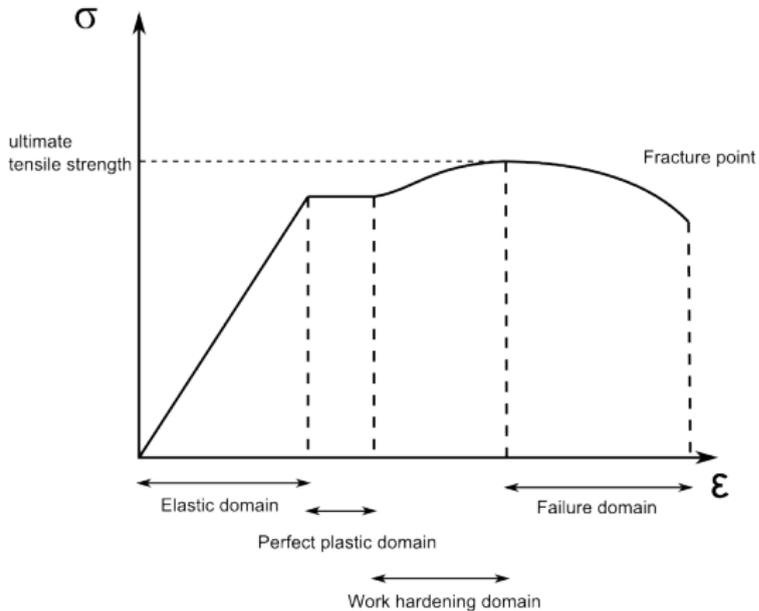
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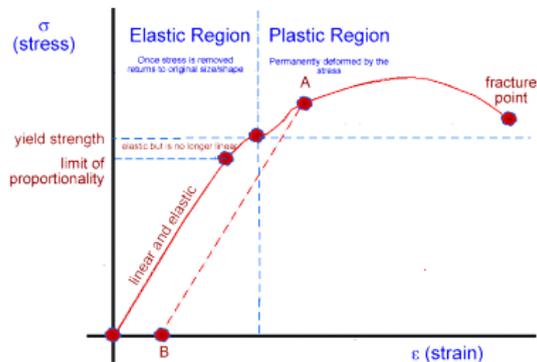


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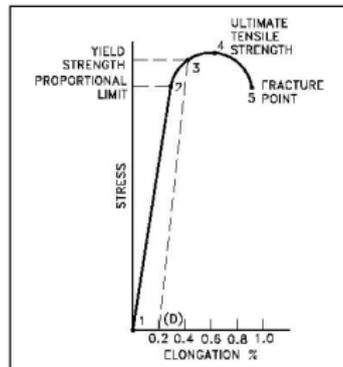
Brittle vs ductile

When the stresses are too large, the system response is not linear anymore

- if the system is brittle, it will break when the stress/strain will reach a critical value
- if the system is ductile, it will sustain an irreversible deformation before finally breaking



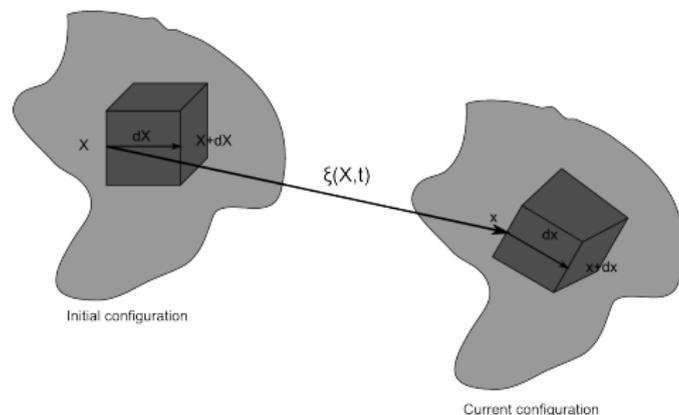
ductile fracture



brittle material

Strains and strain tensor

How to describe more complex solicitations and deformations ?



$$\underline{x} = \underline{X} + \underline{\xi}(X, t)$$

Small deformation hypothesis

Most of the mechanics is made under this approximation. $|\partial \xi_i / \partial X_j| \ll 1$

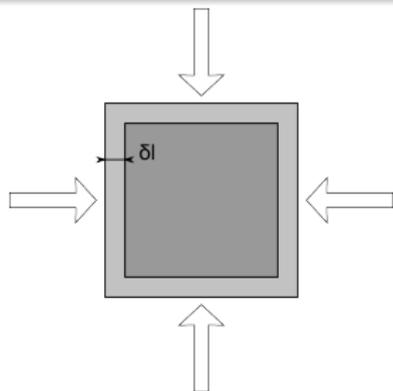
Under this condition, we can define the strain tensor characterizing the change of length and angle of a material vector:

$$\underline{\varepsilon} = \frac{1}{2} (\nabla \underline{\xi} + \nabla^t \underline{\xi}) \quad \varepsilon_{ij} = \frac{1}{2} \left(\frac{\partial \xi_i}{\partial x_j} + \frac{\partial \xi_j}{\partial x_i} \right)$$

Strains and strain tensor

Volumetric strain

$$\varepsilon_{ij} = \frac{1}{2} \left(\frac{\partial \xi_i}{\partial x_j} + \frac{\partial \xi_j}{\partial x_i} \right) \Rightarrow \begin{pmatrix} \varepsilon_{11} & \varepsilon_{12} & \varepsilon_{13} \\ \varepsilon_{21} & \varepsilon_{22} & \varepsilon_{23} \\ \varepsilon_{31} & \varepsilon_{32} & \varepsilon_{33} \end{pmatrix}$$



$$\frac{d\Omega}{\Omega} = 3 \frac{\delta l}{l} = 3\varepsilon = \text{tr}(\underline{\underline{\varepsilon}})$$

$$\underline{\underline{\varepsilon}} = \begin{pmatrix} \varepsilon & 0 & 0 \\ 0 & \varepsilon & 0 \\ 0 & 0 & \varepsilon \end{pmatrix}$$

The diagonal term ε_{ii} of the linearized strain tensor corresponds to a volumetric deformation in the direction e_i .

The change in volume is obtained by $\Omega = (1 + \epsilon) \Omega_0$ with $\epsilon = \text{tr}(\underline{\underline{\varepsilon}})$

Strains and strain tensor

Deviatoric strain

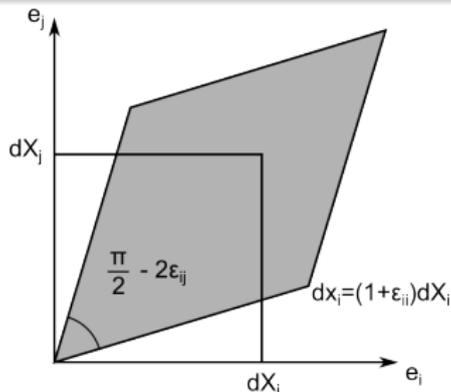
The non diagonal terms are called *deviatoric*

These terms correspond to a change of orientation of a material vector.

The strain tensor is symmetric and its eigenvectors determine the *principal directions of deformation*

Interpretation of non diagonal terms

$2\varepsilon_{ij}$ is the change of angle undergone by two initially normal vectors in the directions i and j .



Decomposition into volumetric and deviatoric strains

$$\underline{\underline{\varepsilon}} = \frac{\epsilon}{3}\underline{\underline{1}} + \underline{\underline{e}}$$

Stresses and stress tensor

Body and surface forces

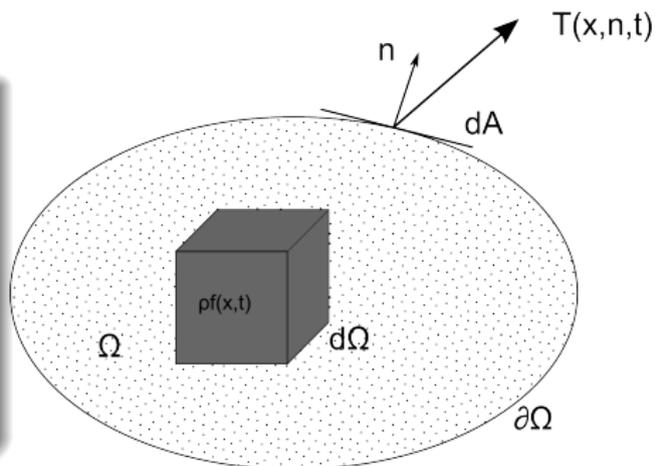
Body forces and contact forces

Body forces are applied on the whole domain Ω

- gravity
- electrostatic field

Contact forces are applied on the boundary of the domain $\partial\Omega$

- pressure
- shear



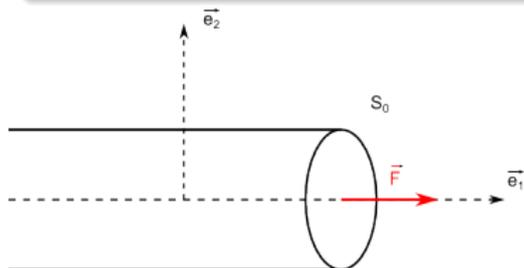
Static equilibrium of a domain

$$\int_{\Omega} \rho \underline{f}(\underline{x}, t) d\Omega + \int_{\partial\Omega} \underline{T}(\underline{x}, t) dA = 0$$

\underline{n} normal to the surface, ρ density of the solid

Stress vector

Stresses are contact forces

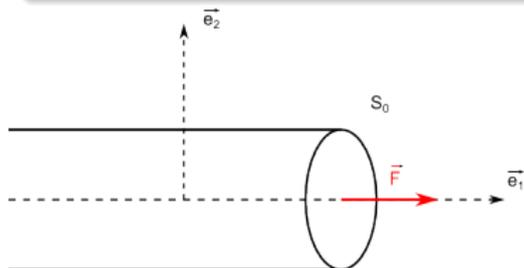


The expression of the stress vector depends on the direction we are looking at:

$$\sigma_0 = \frac{F}{S_0} \quad \tau_0 = 0 \quad \text{and} \quad \underline{T} = \sigma_0 \underline{e}_1 + 0 \times \underline{e}_2$$

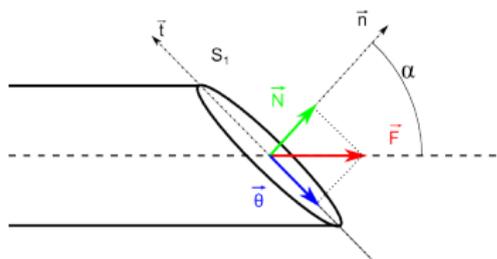
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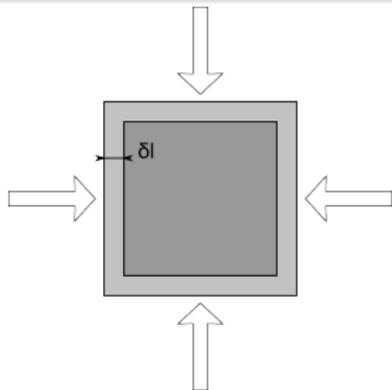
$$\sigma_1 = \frac{F \cos \alpha}{S_1} \quad \tau_1 = \frac{-F \sin \alpha}{S_1} \quad \text{and} \quad \underline{T} = \sigma_1 \underline{n} + \tau_1 \underline{t}$$

Stresses and stress tensor

The contact forces can be expressed as a result of a linear operator in a defined base

$$\underline{T}(\underline{x}, t, \underline{n} = n_j \underline{e}_j) = \underline{\underline{\sigma}} \cdot \underline{n} = \sigma_{ij} n_j \underline{e}_i$$

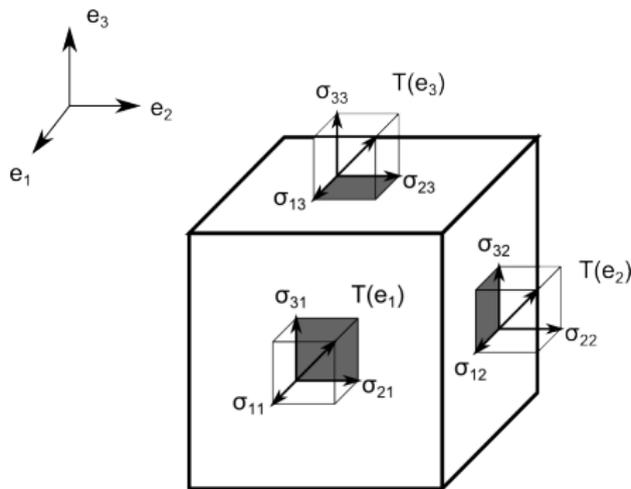
This defines the stress tensor



The stress here is analogous to a pressure. There is a negative sign because the normal to the face is toward the outside

$$\underline{\underline{\sigma}} = \begin{pmatrix} -p & 0 & 0 \\ 0 & -p & 0 \\ 0 & 0 & -p \end{pmatrix}$$

Stresses and stress tensor



Characteristics of the stress tensor

- symmetric tensor. The eigenvector define the principal directions of stress
- σ_{ij} is the stress applying on the surface normal to \underline{e}_i in the direction \underline{e}_j
A pressure corresponds to a negative value
- σ_{ij} is the shear stress acting in along the \underline{e}_j on the surface normal to \underline{e}_i
- the stress tensor can be decomposed as spherical part and deviatoric part
 $\underline{\underline{\sigma}} = \underline{\underline{s}} + \sigma \underline{\underline{1}}$ with $\sigma = 1/3 \text{tr}(\underline{\underline{\sigma}})$
- For a fluid, the shear stress is equal to zero and the spherical part reduces to the pressure with a $-$ sign

Terminology of states of stress

- Hydrostatic pressure : $\sigma_1 = \sigma_2 = \sigma_3$
No shear stress on any plane \rightarrow isotropic pressure
- Uniaxial stress :
 $\sigma_1 > \sigma_2 = \sigma_3 = 0$ (uniaxial traction)
 $\sigma_1 < \sigma_2 = \sigma_3 = 0$ (uniaxial compression)
- Axial stress with confinement : $\sigma_1 > \sigma_2 = \sigma_3 > 0$
- Pure triaxial stress : $\sigma_1 > \sigma_2 > \sigma_3$

Young relation

We can rewrite the Hooke's law with the tensor notation

In the basis of the principal directions of stress, we have the relation:

$$\begin{pmatrix} \varepsilon_1 & 0 & 0 \\ 0 & \varepsilon_2 & 0 \\ 0 & 0 & \varepsilon_3 \end{pmatrix} = \frac{\sigma}{E} \begin{pmatrix} 1 & 0 & 0 \\ 0 & -\nu & 0 \\ 0 & 0 & -\nu \end{pmatrix}$$

In the general case, we have : $\varepsilon_{ij} = \frac{1}{E} \left[(1 + \nu) \sigma_{ij} - \nu \text{tr} \left(\underline{\underline{\sigma}} \right) \delta_{ij} \right]$

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The mechanical equilibrium then writes $\int_{\Omega} \rho \underline{\underline{f}} d\Omega + \int_{\partial\Omega} \underline{\underline{\sigma}} \cdot \underline{\underline{n}} dA = 0$

Which reduces to $\nabla \cdot \underline{\underline{\sigma}} + \rho \underline{\underline{f}} = 0$

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Strain work

The infinitesimal work dW supplied between t and $t + dt$ is:

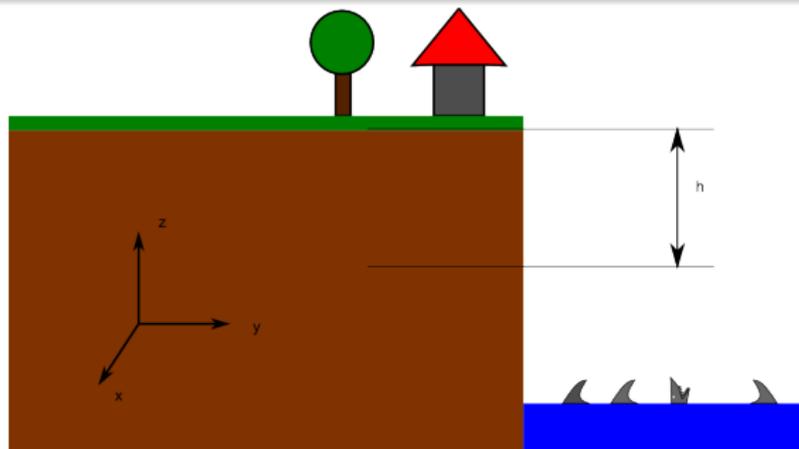
$$dW = \int_{\Omega} d\underline{\underline{\xi}} \cdot \rho \underline{f} d\Omega + \int_{\partial\Omega} (d\underline{\underline{\xi}} \cdot \underline{T}) dA$$

Using the definition of the stress tensor we finally obtain the strain work for infinitesimal transformation:

$$dW = \sum_i \sum_j \sigma_{ij} d\varepsilon_{ij} \text{ which also written as } dW = \underline{\underline{\sigma}} : d\underline{\underline{\varepsilon}}$$

Landslides

Should you really build your house there ?



How can we assess the mechanical equilibrium of a soil ?

Mohr circles

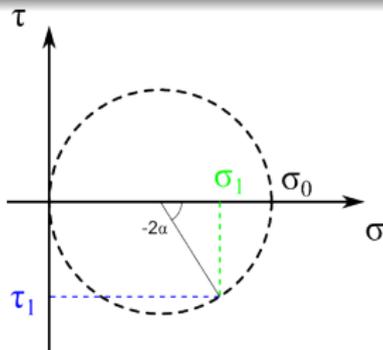
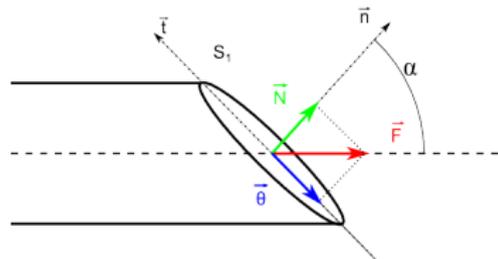
Mohr circles is a way to represent the stresses acting on the material point of a system subjected to a sollicitation

It is a 2D representation of the normal and shear stress with every possible orientation of the cutting plane

Mohr circle for a simple traction test

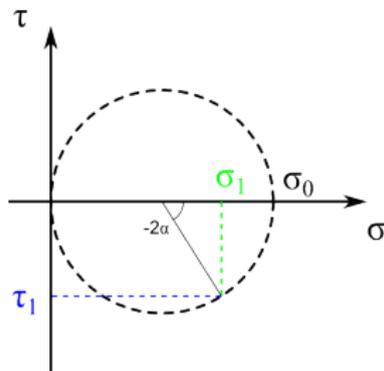
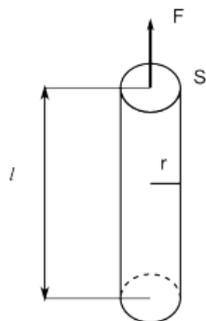
There is only one principal direction of stress: $\underline{\underline{\sigma}} = \begin{pmatrix} \sigma_0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$

$$\begin{cases} \sigma_1(\alpha) = \frac{F \cos \alpha}{S_1(\alpha)} = \sigma_0 \cos^2 \alpha \\ \tau_1(\alpha) = -\frac{F \sin \alpha}{S_1(\alpha)} = -\sigma_0 \cos \alpha \sin \alpha \end{cases} \Rightarrow \text{Parametric equation of a circle}$$



Mohr circles

Failure of a sample on a simple traction test



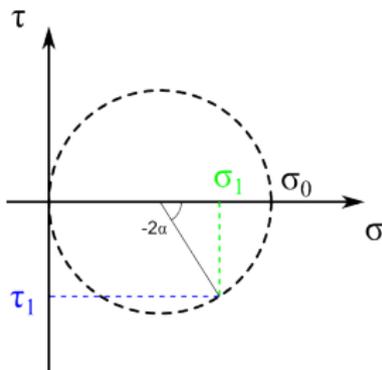
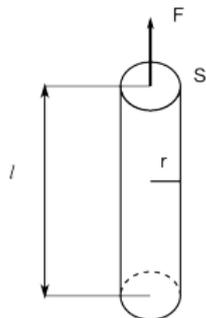
Failure of a ductile material is always due to the shear stress

The maximum shear stress is obtained for an orientation of the plane of 45° from the sollicitation.

The failure will then follow this plane

Mohr circles

Failure of a sample on a simple traction test



Failure of a ductile material is always due to the shear stress

The maximum shear stress is obtained for an orientation of the plane of 45° from the sollicitation.

The failure will then follow this plane

Mohr circles

More complex solicitations

Mohr circles for a biaxial solicitation $\underline{\underline{\sigma}} = \begin{pmatrix} \sigma_1 & 0 & 0 \\ 0 & \sigma_2 & 0 \\ 0 & 0 & 0 \end{pmatrix}$

We choose a cutting plan with an angle α with \underline{e}_1 (\underline{e}_3 is the axis of rotation)
 We can consider without any loss of generality that $\sigma_1 > \sigma_2$

The new normal and tangential components for this plan are:

$$\sigma_n = \frac{F_1 \cos \alpha}{S_1} + \frac{F_2 \sin \alpha}{S_2} = \sigma_1 \cos^2 \alpha + \sigma_2 \sin^2 \alpha$$

$$\tau = -\frac{F_1 \sin \alpha}{S_1} + \frac{F_2 \cos \alpha}{S_2} = -\sigma \sin \alpha \cos \alpha + \sigma_2 \sin \alpha \cos \alpha$$

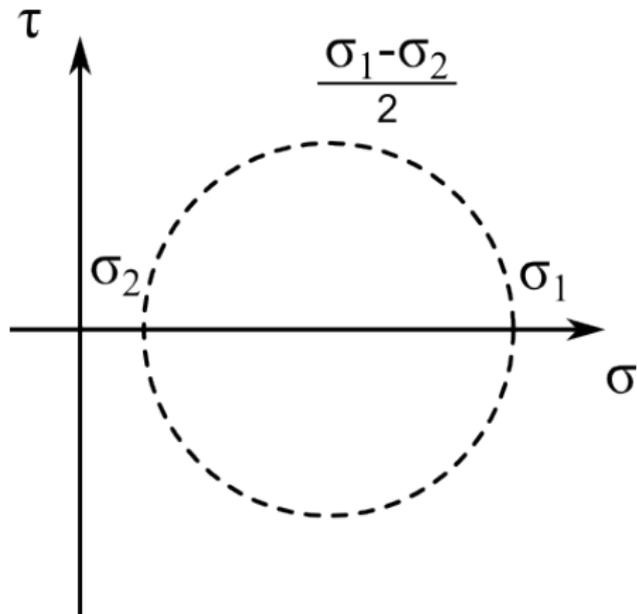
This corresponds to a circle with parametric equation:

$$\sigma_n = \frac{\sigma_1 + \sigma_2}{2} + \frac{\sigma_1 - \sigma_2}{2} \cos(-2\alpha)$$

$$\tau = \frac{\sigma_1 - \sigma_2}{2} \sin(-2\alpha)$$

Mohr circles

More complex solicitations

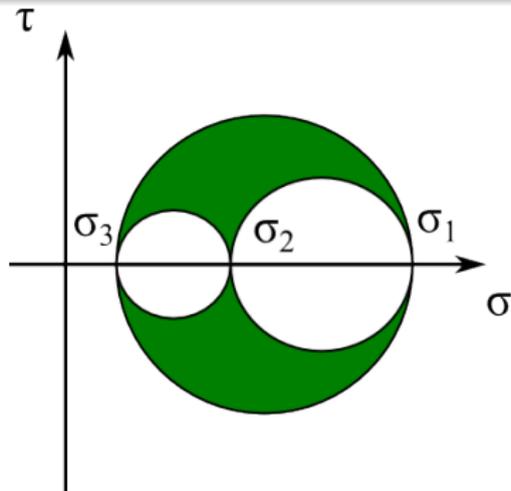


Mohr circles

More complex solicitations

Triaxial sollicitation: $\underline{\underline{\sigma}} = \begin{pmatrix} \sigma_1 & 0 & 0 \\ 0 & \sigma_2 & 0 \\ 0 & 0 & \sigma_3 \end{pmatrix}$

A triaxial sollicitation can be considered as a superposition of 3 biaxial sollicitations with the rotation of the plane around the 3 principal directions of stress. The result is a surface.

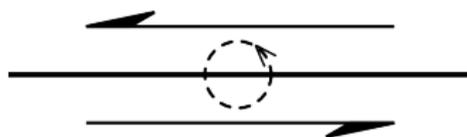


Sign convention in Mohr circles

Mohr circles are usually used in geotechnics: almost all stresses are compression

- compression stresses are positive (opposite from the normal case)
- The stress tensor is antisymmetric

$$\sigma_{12} = -\sigma_{21}$$



Counterclockwise = positive



Clockwise = negative

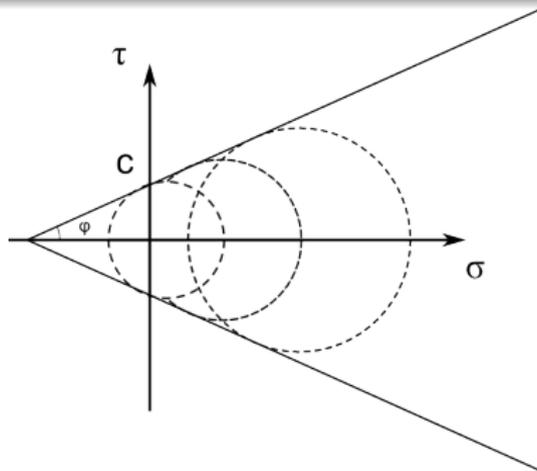
Failure of soils

Mohr-Coulomb criterion for compression failure

$$\tau = \sigma \tan \varphi + c$$

⇒ Failure of a granular material: shear of grains with respect to each other ↔ friction angle φ

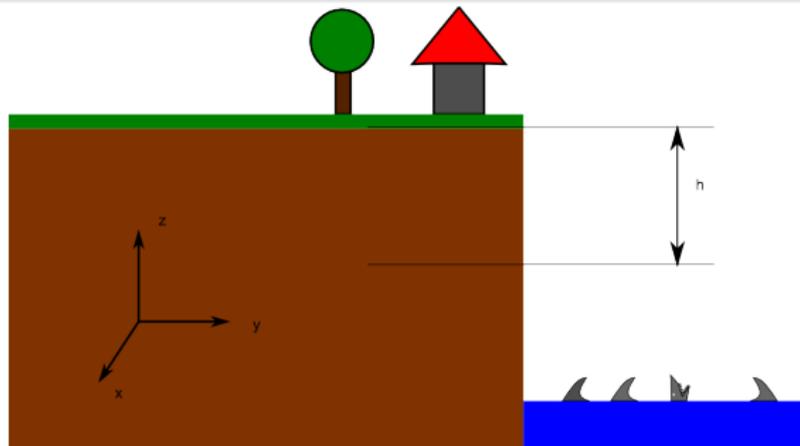
Criterion increasing with normal stress: consolidation



Failure occurs if a Mohr circle reaches the failure envelope. The failure plane orientation is given by the contact point between the circle and the envelope

Landslides

Should you really build your house there ?



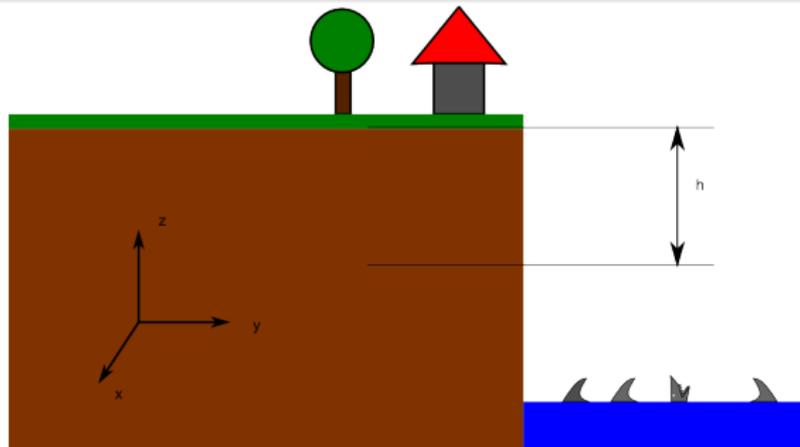
Stability of the soil without house

Vertical stress: $\sigma_{zz} = \rho gh$

$$\underline{\underline{\varepsilon}} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & \varepsilon_{yy} & 0 \\ 0 & 0 & \varepsilon_{zz} \end{pmatrix} \Rightarrow \begin{cases} 0 = \frac{1}{E} (\sigma_{xx} - \nu (\sigma_{yy} + \sigma_{zz})) \\ \varepsilon_{yy} = \frac{1}{E} (\sigma_{yy} - \nu (\sigma_{xx} + \sigma_{zz})) \\ \varepsilon_{zz} = \frac{1}{E} (\sigma_{zz} - \nu (\sigma_{xx} + \sigma_{yy})) \end{cases} \quad \text{Hooke's law}$$

Landslides

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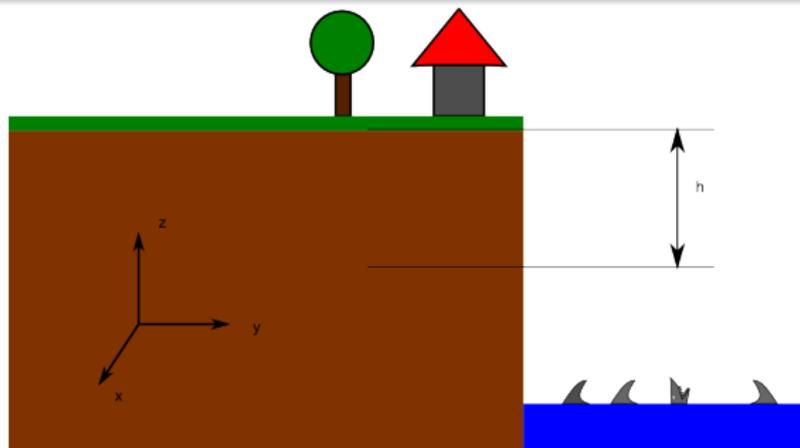


Static equilibrium

$$\nabla \cdot \underline{\underline{\sigma}} - \rho \underline{\underline{g}} h = 0 \Rightarrow \begin{cases} \frac{\partial \sigma_{xx}}{\partial x} = 0 \\ \frac{\partial \sigma_{yy}}{\partial y} = 0 \\ \frac{\partial \sigma_{zz}}{\partial z} = 0 \end{cases} \quad \text{At the cliff, } \sigma_{yy} = 0 \text{ so } \forall y, \sigma_{yy} = 0$$

Landslides

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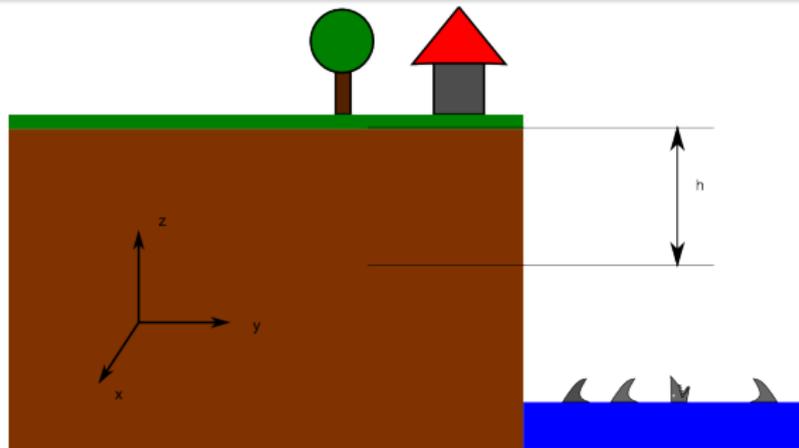


Resolution of the system of equations

$$\begin{cases} \sigma_{xx} = \nu \sigma_{zz} \\ \varepsilon_{yy} = \frac{\nu(1+\nu)}{R_2} \sigma_{zz} \\ \varepsilon_{zz} = \frac{1-\nu^2}{E} \sigma_{zz} \end{cases} \Rightarrow \underline{\underline{\sigma}} = \begin{pmatrix} \nu \rho g h & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & \rho g h \end{pmatrix}$$

Landslides

Should you really build your house there ?

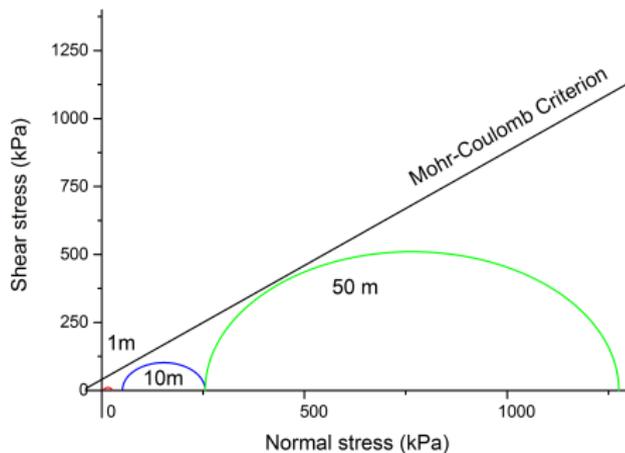
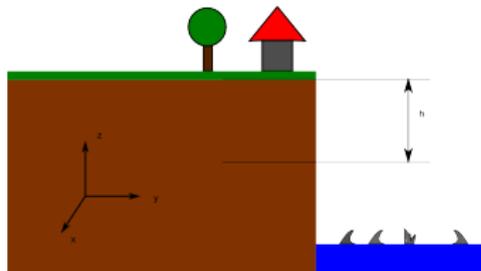


Stability of the soil

$$\left\{ \begin{array}{l} \varphi = 40^\circ \\ c = 40\text{kPa (loam)} \\ \nu = 0.2 \end{array} \right.$$

Landslides

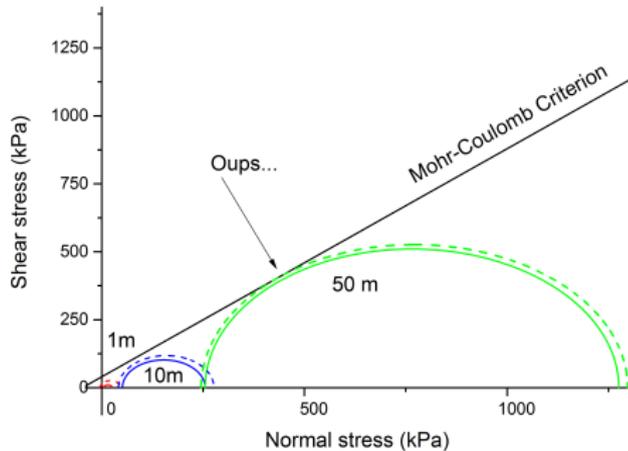
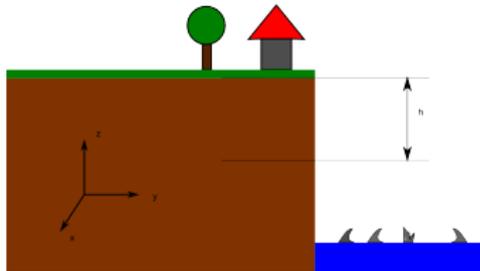
Should you really build your house there ?



Influence of the house: $\sigma_{zz} = \rho gh + P$ (with $M_{house} = 40\text{kPa}$)

Landslides

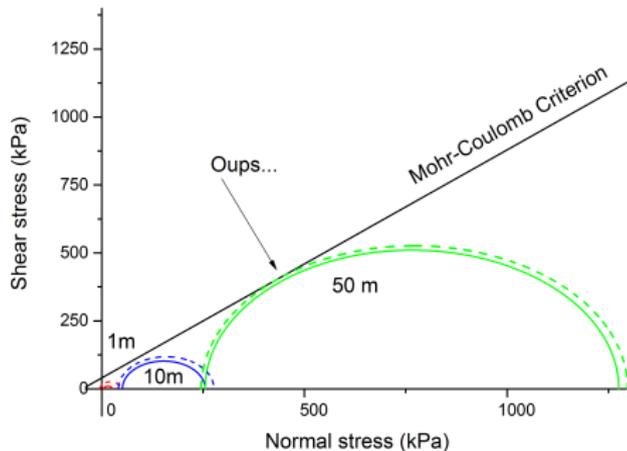
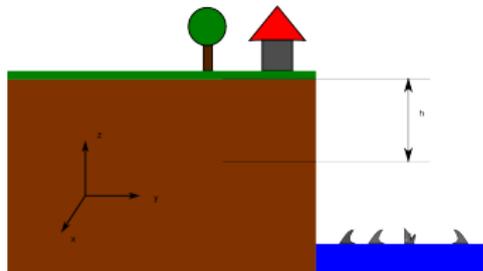
Should you really build your house there ?



Angle of landslide ?

Landslides

Should you really build your house there ?



What of the influence of water in the interstitial space ?

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